

Breather and rogue wave solutions of coupled derivative nonlinear Schrödinger equations

Xiao-Shuo Xiang - Da-Wei Zuo

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Abstract Coupled derivative nonlinear Schrödinger (cDNLS) equations play an important role in plasma physics, optics and other fields. One-order analytical solutions of the cDNLS equations are obtained by virtue of the Darboux transformation. Via adjustment about the parameters, bright-dark conversion mechanism of rogue wave is obtained; combinations of rogue wave and breather/bellshape soliton with different waveform are gotten; distance between rogue wave and breather/bellshape soliton can be changed. In addition, we find that rogue wave has affect on the propagation direction of breather/bellshape soliton.

Keywords Breather and rogue wave solution -Coupled derivative nonlinear Schrödinger equation

1 Introduction

Rogue wave (RW) is a kind of nonlinear wave with wave height more than 2-2.2 times of the initial wave height, which can be seen in the open sea surface but cannot be predicted [1]. RW can be described theoretically by nonlinear evolution equations, and it appears in nonlinear optical systems, water tanks and plasma [2-5]. Periodic RW constitutes breather [6] in which the pulse energy and wave amplitude change periodically during transmission [7,8]. Breather can be explained by nonlinear evolution equation, while it has also been used to describe the nonlinear Alfvén waves in plasma.

When we discuss the slowly varying waves in the plasma, it is known to us that the following assumptions need to be considered: Wave is quasi-sine and is described by the cold plasma equations; plasma consists of the electrons and one kind of ions; all irreversible effects are neglected; dispersive effects come from the electron inertia and Hall term in the generalized Ohm's law. We will consider the waves which propagate in the x direction and assume all quantities depend on x and t only. Base on the above assumptions, controlled equations can be written as [9]

$$\begin{split} \frac{\partial \rho}{\partial t} + \frac{\partial \rho v}{\partial x} &= 0, \\ \frac{\partial \rho v}{\partial x} + \frac{\partial (\rho v^2 + 0.5B^2)}{\partial x} &= 0, \\ \frac{\partial \tilde{u}}{\partial t} + \frac{v \partial \tilde{u}}{\partial x} &= \frac{\partial \tilde{B}}{\rho \partial x}. \\ \frac{\partial \tilde{B}}{\partial t} + \frac{\partial (v \tilde{B} - \tilde{u})}{\partial x} + i \frac{\partial}{\partial x} \left(\frac{\partial \tilde{B}}{\rho \partial x} \right) \\ &- \alpha \frac{\partial}{\partial x} \left(\frac{\partial}{\partial t} + v \frac{\partial}{\partial x} \right) \frac{\partial \tilde{B}}{\rho \partial x} &= 0, \end{split}$$

$$(1)$$

where ρ is the normalized mass density dependent on reality density, v is the normalized fluid velocity in the x direction from $v_A = (\widehat{B}_0^2/4\pi\rho_0)^{1/2}$, while the magnetic field component in the x direction \widehat{B}_0 is a constant by the condition $\nabla \widehat{B} = 0$. We expand ρ , v, \widehat{u} , \widehat{B} and

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Rogue Equations

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