EMITTANCE SCANNER OPTIMIZATION FOR LOW-ENERGY ION BEAMS

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Abstract

Ion beam emittances are normally measured as twodimensional distributions of the beam current fraction within a narrow window Δx centered at position coordinate x and a narrow window $\Delta x'$ centered at trajectory angle x'. The small fraction of the beam current found within both of these windows causes the measured emittance signals to be sensitive to noise, bias, and other unwanted signals. One example of unwanted signals is slit scattering that is discussed in detail for low-energy ion beams.

INTRODUCTION

The emittance is the key parameter to calculate the transport and losses of particle beams. It describes the six-dimensional distribution of all position coordinates along the three configuration space directions and the associate velocity coordinates. It is normally reduced into three subsets by projecting it into the two-dimensional planes [x,x'], [y,y'], and [x,x'].

The experimental projection of transverse subsets, either x or y, is accomplished with a position slit placed at a number of equidistant coordinates. The slit accepts a narrow band Δx or Δy , respectively, of the beam cross section. Further downstream, a second, parallel probe samples narrow bands $\Delta x'$ or $\Delta y'$ of the corresponding trajectory angle distribution x' or y', respectively.

Wire harps can measure each distribution in a single shot, but the accuracy is limited by the uniformity of the amplifier gain, the wires' size, the wires' straightness, and the wires' surface condition that affects the secondary electron emission rate. The wires are likely to be uniform initially, but exposure to the beam will gradually alter some of the characteristics at rates that depend on the exposure, which varies from wire to wire. Uniformity problems can be avoided with a single slit mounted on a suppressed Faraday cup that scans the trajectory angle distribution. Highly accurate and reliable results can be expected if the beam remains stable during the timeconsuming scan.

Measuring the two-dimensional distribution with adequate resolution requires a narrow slit width or thin wires. Hence most of the beam is intercepted on the entrance slit, and only a tiny fraction of the beam is measured. Due to their small amplitude, the signals are sensitive to noise, bias, and signals produced by higher order processes such as slit scattering. The effect of noise and bias can be eliminated with a self-consistent exclusion analysis as discussed in previous papers [1-3]. Data affected by higher order processes can be improved by applying model-based corrections with an accuracy limited by the model. More accurate and reliable data are normally obtained with design modifications that minimize the sensitivity to higher order processes.

Recently we have shown that data from electric sweep scanners contain signals from the beam being dumped on the electrical deflection plates. We have also shown that the amplitude of these ghost signals can be reduced by over 99% when using staircase deflection plates like the ones shown in Fig. 1 [4,5].

In this paper we discuss slit scattering and how it can be minimized for low-energy ion beams. The discussion is valid for most emittance devices scanning low-energy beams. As an example, we use the Allison scanner because it promises the most accurate emittance measurements for low-energy ion beams.

ALLISON SCANNERS

Electric sweep scanners [6] and Allison scanners [7] were introduced more than 25 years ago. Both scanners use an electric field between the two defining slits to measure the distribution of the entrance angle x' of the beamlet passing through the first slit. Allison scanners combine the electrical sweep of the entrance-angles with a mechanical position scan that probes the position distribution x. This allows for mounting both slits on a single support base, enabling relative alignment of both slits within tight tolerances, a potential source of error. Charged particles passing through both slits are collected in a Faraday cup with suppressed secondary electron emission. A grounded light-tight shield surrounds the entire assembly to block unwanted charged particles from reaching the Faraday cup.

The voltage to angle conversion is $x' = V \cdot L_{eff}(2 \cdot g \cdot U)$, where U is the ion's kinetic energy divided by its charge, L_{eff} is the effective length of the deflection plates, g is the

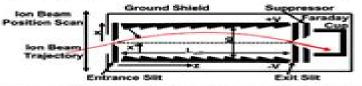


Figure 1: Schematic Allison scanner with staircase deflection plates.

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