

LASERS AND THEIR APPLICATIONS

Passive Mode-Locking of CW Solid-State Lasers with a Semiconductor Saturable Absorber

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Abstract—The conditions of stabilization of ultrashort pulses (USP) in cw solid-state lasers with the saturable semiconductor absorber are analyzed on the basis of the self-consistent field theory and the fluctuation model. It is shown that the nonlinear refractivity can provide the generation of pulses with a duration considerably shorter than the relaxation time of the passive shutter in a wide range of negative dispersion.

INTRODUCTION

In recent years, considerable progress has been achieved in the field of generation of ultrashort pulses (USP) with femtosecond durations. This became possible due to the new approach to mode-locking of solid-state lasers, which is based on the use of phase optical nonlinearities and, first of all, nonlinear refractivity [1]. Latest studies [2] showed that efficient mode-locking in the femtosecond region (up to several femtoseconds) can also be realized by using modulators on the basis of semiconductor structures. Since the relaxation time of the fast component in such structures is not less than 100 fs, and for the slow component it makes up several units or tens of picoseconds, the ordinary passive mode-locking caused by saturation of losses in a semiconductor is inefficient. In this connection, it has been suggested in [2–4] that in this situation, the nonlinear phase effects have a crucial action on the formation of USPs. However, the mechanism of such action was assumed to be different, and, what is quite essential, the specific character of the anomalously high nonlinear refractivity in semiconductors (see [5–7]) was not taken into account.

Semiconductor, such as GaAs, InSb, PbSe, and others, are known to have the anomalously high nonlinearity [5, 7]. In such materials, the coefficient of the dielectric susceptibility depends on the concentration of free carriers and can be different for different levels of excitation, so that the phase mismatch is proportional to the population difference ΔN of the levels (in the approximation of a two-level medium): $i\zeta_2\Delta N$, where ζ_2 is the nonlinear coefficient determined by the difference in the susceptibilities for the ground and excited states [8]. As the result, with the change in the population difference caused by the bleaching of losses, self-phase modulation (SPM) appears, leading to the nonlinear frequency shift and additional amplitude modulation in the presence of the pulse frequency shift with respect to the resonance absorption frequency of the semiconductor [6, 7]. It was shown in [9] that the first of these

effects can stabilize solitons propagating in optical fiber communication lines.

In this paper, we study the conditions of formation and stability of USP in a cw solid-state laser when SPM takes place in the active element and the semiconductor shutter in the presence of the group-velocity dispersion, as well as the energy saturation of losses and gain. We show that due to the existence of the nonlinear refractivity in the semiconductor shutter, the SPM substantially transforms parameters of the USP of the femtosecond duration and leads to its stabilization, despite the fact that the pulse duration is considerably shorter than the excitation relaxation time of the passive shutter.

RESULTS AND DISCUSSION

The analysis is based on the theory of self-consistent field [10]. The equation describing the evolution of the complex field envelope $a(k, t)$ (k is the number of a transit through the cavity, and t is the local time) in the distributed laser system, containing an active medium with the gain α and a semiconductor passive shutter with the absorption coefficient γ , in the presence of dispersion and with allowance for SPM caused by Kerr nonlinearity in the active medium and the nonlinear refractivity in the semiconductor, can be represented in the following form

$$\frac{\partial a(k, t)}{\partial t} = \left[\frac{\alpha - \gamma + i\chi\gamma}{1 + \frac{\partial}{\partial t}} - l + id\frac{\partial^2}{\partial t^2} - i\beta|a(k, t)|^2 + i\Phi \right] a(k, t). \quad (1)$$

Here, the times and the dispersion coefficient d are normalized to the inverse bandwidth Γ of transmission of the frequency filter, χ is the nonlinear refraction coefficient in the semiconductor shutter, β is the coefficient of SPM in the active medium, and l is the linear losses.

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